Spatial distributions of energy and stresses in hadrons

Adam Freese

University of Washington

October 24, 2022

Introduction

- How are energy and momentum spatially distributed in hadrons?
 - Mass decomposition
 - Spin decomposition
 - Distribution of pressures, stresses, forces
 - Energy-momentum tensor (EMT)
- $\bullet\,$ Hadrons are manifestly relativistic.
 - Their size is comparable to their Compton wavelength.
 - Their constituents move close to the speed of light.
 - Light front coordinates provide appropriate description.
- Based on work in:
 - AF & Gerald Miller, PRD103, 094023
 - AF & Gerald Miller, PRD**104**, 014024
 - AF & Gerald Miller, PRD105, 014003
 - AF & Wim Cosyn, arxiv:2207.10787
 - AF & Wim Cosyn, arxiv:2207.10788
 - AF & Gerald Miller, arxiv:2210.03807

Barycenters and relative coordinates

- Say you have two particles bound into one.
 - A hydrogen atom, positronium, ...
- Two kinds of spatial extent:
 - The distance between the bound objects.
 - 2 The spread in the objects' wave functions.
- Barycentric and relative coordinates:

 $\mathbf{R} = m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2$ $\mathbf{r} = \mathbf{r}_2 - \mathbf{r}_1$ $\psi(\mathbf{r}_1, \mathbf{r}_2) = \Psi_{\text{bar}}(\mathbf{R})\psi_{\text{rel}}(\mathbf{r})$

- ${\bf R}$ says where the joint system is. (Position state: $|{\bf R}\rangle).$
- $\bullet~\mathbf{r}_{\mathrm{rel}}$ describes intrinsic structure.
 - Want **only** this.
 - The other piece gives state preparation dependence.



Density as a convolution: non-relativistic case

• Physical densities given by expectation values:

$$\rho_{\rm phys}(\mathbf{x},t) = \langle \Psi | \hat{\rho}_{\rm op.}(\mathbf{x},t) | \Psi \rangle$$

• Galilean symmetry means **barycentric position** and **impact parameter** dependence can be separated.

$$\rho_{\rm phys}(\mathbf{x},t) = \int d^3 \mathbf{R} \left| \Psi_{\rm bar}(\mathbf{R},t) \right|^2 \rho_{\rm internal}(\mathbf{x}-\mathbf{R})$$

• Internal number density:

$$\rho_{\text{internal}}(\mathbf{b}) = \left| \psi_{\text{rel}} \left(\frac{m_1 + m_2}{m_2} \mathbf{b} \right) \right|^2 + \left| \psi_{\text{rel}} \left(-\frac{m_1 + m_2}{m_1} \mathbf{b} \right) \right|^2$$

The problem with barycenters in relativity

- Problem: can't isolate barycentric coordinate in relativistic instant form.
 - *z*-direction boost at time *t*:



- Transformed constituent coordinates at **different times**.
- Simultaneity is relative : boosts transform bound constituents to different times!
- $\bullet\,$ Cannot break density into barycentric spread \otimes internal structure at fixed time.
- Light front coordinates fix this by defining a new, boost-invariant time variable!

Light front coordinates

Light front coordinates are a different foliation of spacetime.



6 / 38

э

Light front: Myths and Facts

- Myth: Light front coordinates are a reference frame.
- Fact: Light front coordinates can be employed in any reference frame.
- Myth: Light front coordinates describe the perspective of light.
- Fact: Light front coordinates describe *our* perspective.
 - ... but only in the z direction!
 - Light has no perspective.
- Myth: Light front coordinates come from boosting to infinite momentum.
- Fact: Light front coordinates come from redefining:
 - Interval and the second sec
 - **2** What we mean by *boosting*
 - (a) How we break the Poincaré group into generators (Galilean subgroup)

Light front & time synchronization

Light front redefines simultaneity

- Fixed $x^+ = \frac{t+z}{\sqrt{2}}$ means simultaneous
- Look in the $+\hat{z}$ direction...
 - Whatever you **see** right now, is **happening** right now.
 - Only true for $+\hat{z}$ direction though.
- Light front coordinates are what we see.
 - ...at least in one fixed direction.
 - (great for small systems—hadrons!)
 - Not what light "sees."



Terrell rotations

- Lorentz-boosted objects appear *rotated*.
 - Terrell rotation
 - Optical effect: contraction + delay
- Light front transverse boost *undoes* Terrell rotation:

$$B_x^{(\mathrm{LF})} = \frac{1}{\sqrt{2}} \Big(K_x + J_y \Big)$$

- Combination of ordinary boost + rotation!
- Leaves x^+ (time) invariant!
- Changes p_z , but leaves p^+ invariant:

$$P^{+} = \frac{E + p_{z}}{\sqrt{2}} = \frac{\sqrt{p_{z}^{2} + \mathbf{p}_{\perp}^{2} + M^{2}} + p_{z}}{\sqrt{2}}$$

• Dice images by Ute Kraus, https://www.spacetimetravel.org/



Galilean subgroup

- Poincaré group has a (2 + 1)D Galilean subgroup.
 - x^+ is time and \mathbf{x}_{\perp} is space under this subgroup.
 - x^- can be integrated out.
 - $P^+ = \frac{1}{\sqrt{2}}(E_{\mathbf{p}} + p_z)$ is the central charge.
 - x^+ and P^+ are invariant under this subgroup!
- Basically, light front coordinates should give a **fully relativistic** 2D picture that looks like *non-relativistic* physics.
 - But with P^+ in place of M. $\frac{\mathrm{d}\mathbf{P}_{\perp}}{\mathrm{d}x^+} = P^+ \frac{\mathrm{d}^2 \mathbf{x}_{\perp}}{\mathrm{d}x^{+2}}$ $H = P^- = H_{\mathrm{rest}} + \frac{\mathbf{P}_{\perp}^2}{2P^+}$ $\mathbf{v}_{\perp} = \frac{\mathbf{P}_{\perp}}{P^+}$

The cost: lose one spatial dimension (2D densities).



Galilean group and densities

• Physical densities given by expectation values:

$$\rho_{\rm phys}(x^+, \mathbf{x}_\perp) = \int \mathrm{d}x^- \langle \Psi | \hat{\rho}_{\rm op.}(x) | \Psi \rangle$$

- Galilean subgroup means **barycentric position** and **impact parameter** dependence can be separated.
 - Allows wave packet dependence to be factored out.
 - Instant form coordinates **do not** have this luxury!
- Simple densities:

$$\rho_{\rm phys}(x^+, \mathbf{x}_{\perp}) = \int d^2 \mathbf{R}_{\perp} \left| \Psi(\mathbf{R}_{\perp}, x^+) \right|^2 \rho_{\rm internal}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp})$$

• Compound densities:

$$\rho_{\rm phys}(x^+, \mathbf{x}_{\perp}) = \int d^2 \mathbf{R}_{\perp} \left\{ S_{\Psi}^{(1)}(\mathbf{R}_{\perp}, x^+) \rho_{\rm int.}^{(1)}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp}) + S_{\Psi}^{(2)}(\mathbf{R}_{\perp}, x^+) \rho_{\rm int.}^{(2)}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp}) \right\}$$

The energy-momentum tensor

- The **energy-momentum tensor** (EMT) is an operator characterizing the distribution and flow of energy and momentum.
 - Source of gravitation
 - Related to spacetime translation symmetry (Noether's theorems)
- Physical densities given by expectation values:

$$T^{\mu\nu}(x^+, \mathbf{x}_\perp) = \int \mathrm{d}x^- \langle \Psi | \hat{T}^{\mu\nu}(x) | \Psi \rangle$$

- 2D density on transverse plane.
- Every component is a simple or compound density.
- Example: P^+ density is a **simple density**:

$$T^{++}(x^+, \mathbf{x}_\perp) = \int \mathrm{d}^2 \mathbf{R}_\perp \left| \Psi(\mathbf{R}_\perp, x^+) \right|^2 T^{++}_{\mathrm{intrinsic}}(\mathbf{x}_\perp - \mathbf{R}_\perp)$$

Form factors of the EMT

- EMT matrix elements give gravitational form factors (GFFs).
 - It's just a name.
 - EMT is the source of gravitation: $G^{\mu\nu} + \Lambda g^{\mu\nu} = 8\pi T^{\mu\nu}$
 - But we don't really use gravitation to measure them.
- Analogy to electromagnetic form factors.
- Spin-zero example:

$$\langle p'|\hat{J}^{\mu}(0)|p\rangle = 2P^{\mu}F(t)$$

$$\langle p'|\hat{T}^{\mu\nu}(0)|p\rangle = 2P^{\mu}P^{\nu}A(t) + \frac{1}{2}(\Delta^{\mu}\Delta^{\nu} - \Delta^{2}g^{\mu\nu})D(t)$$

- A(t) encodes momentum density (mass in NR limit)
- D(t) encodes stress distributions (anisotropic pressures)

$$P = \frac{1}{2}(p+p')$$
$$\Delta = (p'-p)$$
$$t = \Delta^2$$

How to get the GFFs

- Hard exclusive reactions are used to measure GFFs—not gravity experiments.
 - Deeply virtual Compton scattering (DVCS) to probe quark structure.
 - Deeply virtual meson production (DVMP), e.g., J/ψ or Υ to probe gluon structure.
 - ...and more!
- Measured at Jefferson Lab and the upcoming Electron Ion Collider.



Components of the EMT

$$T^{\mu\nu}(x) = \begin{bmatrix} T^{++}(x) & T^{+1}(x) & T^{+2}(x) & T^{+-}(x) \\ T^{1+}(x) & T^{11}(x) & T^{12}(x) & T^{1-}(x) \\ T^{2+}(x) & T^{21}(x) & T^{22}(x) & T^{2-}(x) \\ T^{-+}(x) & T^{-1}(x) & T^{-2}(x) & T^{--}(x) \end{bmatrix}$$

- Momentum densities
- Energy density
- Stress tensor
- $\bullet\,$... and some other things
- Angular momentum density (z component) accessible too:

$$J_z(x) = x^1 T^{+2}(x) - x^2 T^{+1}(x)$$

...basically, from $\mathbf{x}\times\mathbf{p}$

Light front momentum densities

• **Physical** P^+ density is a simple density:

$$T_{\rm phys}^{++}(\mathbf{x}_{\perp}) = P^{+} \int \mathrm{d}^{2}\mathbf{R}_{\perp} \left| \Psi(\mathbf{R}_{\perp}, x^{+}) \right|^{2} T_{\rm int.}^{++}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp})$$

• Transverse momentum density involves *same internal density*:

$$\mathbf{T}_{\mathrm{phys}}^{+i}(\mathbf{x}_{\perp}) = \int \mathrm{d}^{2}\mathbf{R}_{\perp}\Psi^{*}(\mathbf{R}_{\perp}, x^{+}) \frac{-i\overleftarrow{\nabla}_{\perp}^{i}}{2}\Psi(\mathbf{R}_{\perp}, x^{+})T_{\mathrm{int.}}^{++}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp})$$

- But a different smearing function
- Internal density is simple Fourier transform:

$$T_{\rm int.}^{++}(\mathbf{b}_{\perp}) = \int \frac{\mathrm{d}^2 \mathbf{\Delta}_{\perp}}{(2\pi)^2} \frac{\langle p', \mathbf{s}_{\perp} | \hat{T}^{++}(0) | p, \mathbf{s}_{\perp} \rangle}{2(P^+)^2} e^{-i\mathbf{\Delta}_{\perp} \cdot \mathbf{b}_{\perp}}$$

• Internal structure has polarization dependence.

Light front momentum density

• P^+ density is a 2D Fourier transform:

$$\rho_{P^+}^{(\mathrm{LF})}(\mathbf{b}_{\perp}, \mathbf{s}_{\perp}) = \int \frac{\mathrm{d}^2 \Delta_{\perp}}{(2\pi)^2} \frac{\langle p', \mathbf{s}_{\perp} | T^{++}(\mathbf{0}) | p, \mathbf{s}_{\perp} \rangle}{2(P^+)^2} e^{-i\Delta_{\perp} \cdot \mathbf{b}_{\perp}}$$

$$\stackrel{\text{Helicity state}}{\overset{\text{Helicity state}}{\overset{\text{$$

0

0

Spin-one targets Helicity +1

Helicity 0



 P^+ density depends on helicity for spin-one targets. AF & Wim Cosyn, arxiv:2207.10787

Transverse polarization Transverse, $m_s = +1$



Transverse polarization contains helicity-flip contributions. AF & Wim Cosyn, arxiv:2207.10787

A. Freese (UW)

Transverse, $m_s = 0$

Energy density

• Light front energy density is a *compound density*:

$$\begin{aligned} T_{\rm phys}^{+-}(\mathbf{x}_{\perp}) &= \frac{1}{2P^+} \int \mathrm{d}^2 \mathbf{R}_{\perp} \Big\{ \Big| \Psi(\mathbf{R}_{\perp}, x^+) \Big|^2 T_{\rm int.}^{+-}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp}) \\ &- \Psi^*(\mathbf{R}_{\perp}, x^+) \frac{\overleftarrow{\nabla}_{\perp}^2}{4} \Psi(\mathbf{R}_{\perp}, x^+) T_{\rm int.}^{++}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp}) \Big\} \end{aligned}$$

- First piece is true intrinsic energy density
 - Quark mass energy
 - Quark kinetic energy (relative to barycenter)
 - Potential energy & internal stresses
 - Literally the density of the $2P^+P^- \mathbf{P}_{\perp}^2$ operator used by light front folks!
- Second piece is **barycentric kinetic energy**
 - It's literally just the $\mathbf{P}_{\perp}^2/(2P^+)$ density
 - Tells us nothing about internal dynamics
 - Galilean subgroup allows us to isolate it

Example densities from holographic model



Using soft wall holographic model of Brodsky & de Teramond, PRD77 (2008) 056007

Stress

• Stress tensor is also a *compound density*:

$$\begin{split} T^{ij}_{\rm phys}(\mathbf{x}_{\perp}) &= \frac{1}{P^+} \int \mathrm{d}^2 \mathbf{R}_{\perp} \Big\{ \Big| \Psi(\mathbf{R}_{\perp}, x^+) \Big|^2 T^{ij}_{\rm int.}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp}) \\ &- \Psi^*(\mathbf{R}_{\perp}, x^+) \frac{\overleftarrow{\nabla}^i_{\perp} \overleftarrow{\nabla}^j_{\perp}}{4} \Psi(\mathbf{R}_{\perp}, x^+) T^{++}_{\rm int.}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp}) \Big\} \end{split}$$

- First piece is true intrinsic stress tensor
 - Stresses seen by comoving observer
 - Static pressures
- Second piece is stresses from hadron flow
 - Includes motion of hadron
 - Includes wave function dispersion
- Sum of both gives dynamic pressures

Stress tensor and hadron flow

- Compound form of stress tensor mimics classical continuum mechanics
- In Galilean theory (eg **light front**):

 $T^{ij}(\mathbf{x}, \mathbf{v}, \nabla \mathbf{v}) = \mathbf{v}^i \mathbf{v}^j \rho(\mathbf{x}) + T^{ij}_{\text{pure}}(\mathbf{x}, \nabla \mathbf{v})$

- $\mathbf{v}^i \mathbf{v}^j \rho(\mathbf{x})$ depends on wave packet.
- Comoving stress tensor:

$$S^{ij}(\mathbf{x}, \nabla \mathbf{v}) = T^{ij}(\mathbf{x}, \mathbf{v} = 0, \nabla \mathbf{v})$$

- Stresses seen by comoving obsever
- True internal structure of hadron



Example: spin-zero with form factors

• Full stress tensor:

$$\frac{1}{2P^+} \langle p' | \hat{T}^{ij}(0) | p \rangle = P^+ \frac{\mathbf{P}_{\perp}^i}{P^+} \frac{\mathbf{P}_{\perp}^j}{P^+} A(t) + \frac{1}{4P^+} \Big(\mathbf{\Delta}_{\perp}^i \mathbf{\Delta}_{\perp}^j - \mathbf{\Delta}_{\perp}^2 \delta^{ij} \Big) D(t)$$

• Hadron flow:

$$V_{\rm LF}^{ij}(\mathbf{b}_{\perp}) = \left\langle \frac{\mathbf{P}_{\perp}^{i}}{P^{+}} \frac{\mathbf{P}_{\perp}^{j}}{P^{+}} \right\rangle P^{+} \int \frac{\mathrm{d}^{2} \mathbf{\Delta}_{\perp}}{(2\pi)^{2}} A(t) e^{-i\mathbf{\Delta}_{\perp} \cdot \mathbf{b}_{\perp}}$$

• Pure stress tensor:

$$S_{\rm LF}^{ij}(\mathbf{b}_{\perp}) = \frac{1}{4P^+} \int \frac{\mathrm{d}^2 \mathbf{\Delta}_{\perp}}{(2\pi)^2} \Big(\mathbf{\Delta}_{\perp}^i \mathbf{\Delta}_{\perp}^j - \delta^{ij} \mathbf{\Delta}_{\perp}^2 \Big) D(t) e^{-i\mathbf{\Delta}_{\perp} \cdot \mathbf{b}_{\perp}}$$

• Only **D-term** appears in internal stresses.

Comoving stress tensor, D-term, & intrinsic stresses

• General spins: comoving stress tensor *defines* effective **D-term**.

$$S_{\rm LF}^{ij}(\mathbf{b}_{\perp},\mathbf{s}_{\perp}) = \frac{1}{4P^+} \int \frac{\mathrm{d}^2 \mathbf{\Delta}_{\perp}}{(2\pi)^2} \Big(\mathbf{\Delta}_{\perp}^i \mathbf{\Delta}_{\perp}^j - \delta^{ij} \mathbf{\Delta}_{\perp}^2 \Big) D_{\rm eff}(\mathbf{\Delta}_{\perp},\mathbf{s}_{\perp}) e^{-i\mathbf{\Delta}_{\perp}\cdot\mathbf{b}_{\perp}}$$

• D-term encodes intrinsic stresses, or comoving stresses.

• Stresses as seen by observer moving with flow (like a leaf on a river)



Pressures and eigenpressures

• Pressures from **matrix elements**:

 $p_{\hat{n}}(b_{\perp}) = \hat{n}_i \hat{n}_j S_{\rm LF}^{ij}(\mathbf{b}_{\perp})$

- \hat{n} : normal to pressure gauge
- Two **eigenpressures**:

 $S_{\rm LF}^{ij}(\mathbf{b}_{\perp})\hat{n}_j = \lambda_n(b_{\perp})\hat{n}_i$

- Pressure is isotropic *if and only if* eigenpressures are degenerate!
- Anisotropic pressures in general.
- Meaning of sign?
 - **Positive**: gauge is pushed from both sides.
 - **Negative**: gauge is pulled from both sides.



Proton's eigenpressures

Tangential



- Helicity proton (spin in \hat{z} direction)
- Tripole model with f_2 and σ poles
- Streses that are eigenvalues of stress tensor
- **Positive radial pressure** related to D(t) < 0.
- **Polyakov's conjecture**: D(0) < 0 is necessary for stability.

Momentum conservation and force balance

• Conservation law from Noether's theorem:

$$\partial_{\mu}T^{\mu\nu}(x) = 0$$

• Additional **force balance** equation:

$$\mathbf{F}_{\perp}^{j}(\mathbf{x}) = \nabla_{\perp i} S^{ij}(\mathbf{x}) = 0$$



- Force density acting on a hadron is everywhere zero.
 - The hadron is in equilibrium.
 - The hadron is not being acted on by outside forces.
 - Pressure plots are not net force plots!

Pion pressures and energy density

• Stress tensor and energy density are compound densities.

• Internal (pure) densities given by 2D Fourier transforms:

$$\begin{split} T^{ij}_{\text{pure}}(\mathbf{b}_{\perp}) &= \int \frac{\mathrm{d}^2 \mathbf{\Delta}_{\perp}}{(2\pi)^2} \frac{\mathbf{\Delta}_{\perp}^i \mathbf{\Delta}_{\perp}^j - \delta^{ij} \mathbf{\Delta}_{\perp}^2}{2} D(-\mathbf{\Delta}_{\perp}^2) e^{-i\mathbf{\Delta}_{\perp} \cdot \mathbf{b}_{\perp}} \\ \mathcal{E}(\mathbf{b}_{\perp}) &= \int \frac{\mathrm{d}^2 \mathbf{\Delta}_{\perp}}{(2\pi)^2} \left(m_{\pi}^2 - \frac{\mathbf{\Delta}_{\perp}^2}{4} \right) A(-\mathbf{\Delta}_{\perp}^2) e^{-i\mathbf{\Delta}_{\perp} \cdot \mathbf{b}_{\perp}} - \delta_{ij} T^{ij}_{\text{pure}}(\mathbf{b}_{\perp}) \end{split}$$

• Phenomenological form factors:

$$A(t) = \frac{1}{1 - t/m_{f_2}^2} \qquad m_{f_2} = 1270 \text{ MeV}$$
$$D(t) = \frac{-1}{(1 - t/m_{f_2}^2)(1 - t/m_{\sigma}^2)} \qquad m_{\sigma} = 630 \text{ MeV}$$

- Forms inspired by Masjuan et al [PRD87 (2013) 014005]
- Poles chosen to match Kumano's radii [PRD97 (2018) 014020]
- AF & Gerald Miller arxiv:2210.03807 for more info!



Simple densities and wave packet localization

• Simple densities:

$$\rho_{\rm phys}(x^+, \mathbf{x}_{\perp}) = \int d^2 \mathbf{R}_{\perp} \left| \Psi(\mathbf{R}_{\perp}, x^+) \right|^2 \rho_{\rm internal}(\mathbf{x}_{\perp} - \mathbf{R}_{\perp})$$

• Localization at $x^+ = 0$ means:

$$\left|\Psi(\mathbf{R}_{\perp},0)\right|^2 \to \delta^{(2)}(\mathbf{R}_{\perp}), \qquad \qquad \rho_{\rm phys}(0,\mathbf{x}_{\perp}) \to \rho_{\rm internal}(\mathbf{x}_{\perp})$$

- Common practice has thus been to localize to isolate internal densities.
 - $\bullet\,$ Burkardt & Diehl pioneered this idea for light front densities

Int. J. Mod. Phys. A18 (2003) 173, Eur.Phys.J.C 25 (2002) 223-232

- Works in the cases they considered ...
- ...but fails for compound densities.



Pion energy and wave packet localization Internal energy Physical energy (packet dependent)

- Energy density is a compound density.
- Localized packets do not give internal density in this case.
- Packet localization only works for simple densities.



- Sum of five diagrams.
 - Purple lines are gravitons
- Triangle diagrams give spatial extent.
- Gauge invariance imposes:

$$A(t) = D(t) = 1 + \frac{\alpha}{2\pi} A_{\rm LO}(t) + \dots$$

- Singular density at photon center!
 - Density goes as b_{\perp}^{-2} near $b_{\perp} \sim 0$.

Photon target: P^+ density Circularly polarized

Horizontally polarized



- P^+ density in a QED photon (can become virtual e^-e^+)
- Light front densities can describe massless targets!

Photon target: P^+ density (log-scaled) Circularly polarized

Horizontally polarized



- P^+ density in a QED photon (can become virtual e^-e^+)
- Radius 6.5 fm, quadrupole moment 7.7 fm².

Photon target: normal stresses

- QED photon (made of electron & positron)
- Gauge invariance requires that:

A(t) = D(t)

- A(0) = 1 (momentum sum rule)
- This means D(0) = 1 > 0.
- Violates D(0) < 0 condition!
- But photon is stable...

• Radial eigenpressure is negative

- Consequence of D(t) > 0.
- $\bullet\,$ Intrinsic pressure, ${\bf not}$ radiation pressure
- $\bullet\,$ Not a net force plot—net force is zero
- e^-e^+ pulled towards & away from center
- Tangential eigenpressure is (mostly) positive



Horizontally polarized photon



- Eigenpressures & eigenvectors plotted above.
- Formalism has no trouble with anisotropic systems!

Intrinsic pressure: felt within photon.

- Eigenvalues of S^{ij}
- D term only!
- Finite for localized wave packet.
- Can be negative

Radiation pressure: exerted by photon.

- Eigenvalues of T^{ij}
- Combines A and D terms!
- Diverges for localized wave packet.
- Always positive



Summary

- Relativistic densities can be given in **light front coordinates**
- The **energy momentum tensor** encodes momentum, energy & stress distributions.
- Several of these are **compound densities**.
- Energy density decompses into **barycentric kinetic energy** and **internal energy**.
- Stress tensor decomposes into **flow** and **intrinsic stresses**.
 - Flow piece is part of radiation pressure.
- Galilean invariance is needed to allow these decompositions.

Thank you for your time!